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To cite this article: Junli Wang et al 2017 J. Opt. 19 095506

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J. Opt. 19 (2017) 095506 (6pp)

## A high-energy passively *Q*-switched Yb-doped fiber laser based on WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> saturable absorbers

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Received 14 April 2017, revised 26 June 2017 Accepted for publication 12 July 2017 Published 23 August 2017

#### Abstract

In this paper, we report two different saturable absorbers based on WS<sub>2</sub> film and Bi<sub>2</sub>Te<sub>3</sub> film with similar preparation processes. A high-energy stable Q-switching pulse is achieved in identical cavity configurations for each absorber. A modulation depth of 10.17% and a maximum single pulse energy of 56.50 nJ are obtained when employing the WS<sub>2</sub> film. However, using the Bi<sub>2</sub>Te<sub>3</sub> film we obtain a higher modulation depth of 23.04% and a larger single pulse energy of 61.80 nJ, and stable dual-wavelength Q-switching operation was shown at a pump power of 92 mW.

Keywords: WS<sub>2</sub>, Bi<sub>2</sub>Te<sub>3</sub>, Q-switched laser, fiber laser

(Some figures may appear in colour only in the online journal)

#### 1. Introduction

Passively Q-switched fiber lasers have many significant potential applications, for example in remote sensing, medicine, industry and fiber communications [1–3]. Methods of achieving a Q-switched fiber laser include active Q-switched operation and passive Q-switched operation. Compared with actively Q-switched fiber lasers, passively Q-switched lasers are simple, compact and flexible [4, 5]. Passive Q-switched operation with saturable absorbers (SAs) is widely used in pulse fiber lasers. Various types of SAs have been extensively studied and developed in the past few decades, including semiconductor saturable absorber mirrors (SESAMs) [6], single-wall carbon nanotubes (SWCNTs) [7] and graphene [8]. Graphene is the most attractive of these saturable absorbers due to its unique optoelectronic characteristics and broadband saturable absorption effect, but it lack of a bandgap limits its applications



TMDs have unique electronic and optical properties, such as indirect-to-direct bandgap transition with reducing number of layers [15]. TIs have insulating bulky states with an indirect bandgap of 0.35 eV and gapless surface states [16]. Due to their unique characteristics, they have been studied in detail. There are several examples where TMDs and TIs have been used to generate passively *Q*-switched and mode-locked fiber lasers. In 2013, Wang *et al* observed saturable absorption of MoS<sub>2</sub> [17]. Luo *et al* were the first to demonstrate MoS<sub>2</sub>-based *Q*-switched fiber lasers in the wavelength range  $1-2 \mu m$  [18]. Then, Lin *et al* reported a WS<sub>2</sub>-based passively *Q*-switched fiber laser with a wide tunable range at 1.0  $\mu m$  [19]. In 2012, there was the



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Figure 1. Photographs of WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> films.

first demonstration of SA behavior of TIs with a  $Bi_2Te_3$ -based mode-locked fiber laser near 1550 nm [20]. Ever since that discovery,  $Bi_2Te_3$  has been deployed in passively *Q*-switched fiber lasers with different operating wavelengths [21, 22].

In this paper, we demonstrate high-energy Q-switching operation based on WS<sub>2</sub> film and Bi<sub>2</sub>Te<sub>3</sub> film in a Yb-doped ring cavity. The WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> films are fabricated by the same simple method. A nonlinear saturable absorption experiment is carried out to measure the modulation depth of the SAs. The WS<sub>2</sub> film shows a modulation depth of 10.17%, and the Bi<sub>2</sub>Te<sub>3</sub> film exhibits a higher modulation depth of 23.04%. We construct a high-energy Yb-doped Q-switched fiber laser and the results for different materials are compared based on WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> films. Analysis of the different Q-switched performances provides guidance for selecting a suitable SA to satisfy the specific requirements of a Q-switched fiber laser.

#### 2. Material preparation and characterization

A SA in thin-film form provides many advantages for photonic applications, such as flexibility of usage, uniform quality and mass preparation. In our experiment, WS2-polyvinyl alcohol (PVA) and Bi<sub>2</sub>Te<sub>3</sub>-PVA are fabricated by a similar process. WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> dispersions are first fabricated by a liquidphase exfoliation method. The WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> powders were each dispersed in deionized water, using sodium dodecyl sulfate as a surfactant. The dispersions were sonicated for 6 h, then centrifuged at 3000 rpm for 90 min. The top of the dispersions was collected by pipette. Then, the PVA power was dissolved in deionized water with ultrasonic agitation at 90 °C for 3 h. In the next step, the  $WS_2/Bi_2Te_3$  dispersions were mixed with PVA solution and poured into polystyrene cells. Finally, these cells were put into an oven for evaporation at 40 °C. When the cells were dried, WS2-PVA/Bi2Te3-PVA thin films were achieved; these can be used as SAs for Q-switching and mode locking. Photographs of the WS2 and Bi2Te3 films are shown in figure 1.

A Raman spectroscopy system with an excitation wavelength of 633 nm was used to confirm the existence of WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> nanosheets, as shown in figure 1. Figure 2(a) shows the in-plane vibrational mode  $E^{l}_{2g}$  at 350.8 cm<sup>-1</sup> and the out-of-

plane vibrational mode  $A_{1g}$  of WS<sub>2</sub> at 420.2 cm<sup>-1</sup>. Three typical Raman optical photon peaks of Bi<sub>2</sub>Te<sub>3</sub> are shown in figure 2(b), which are identified as  $A_{1g}^{1}$  at 61 cm<sup>-1</sup>,  $E_{g}^{2}$  at 101 cm<sup>-1</sup> and  $A_{1g}^{2}$  at 133 cm<sup>-1</sup>.

Nonlinear saturable absorption is the key parameter for evaluating a SA. We built a twin-detector measurement system to investigate the saturable absorption of  $WS_2$ -PVA and Bi<sub>2</sub>Te<sub>3</sub>-PVA thin films. In the system, a stable home-made picosecond pulsed laser source (central wavelength 1029.8 nm, pulse duration 4.1 ps, repetition rate 27.22 MHz) worked as the illuminator. The measurement results are shown in figure 3. The fitting formula is [23]

$$T(I) = 1 - \Delta T \times \exp(-I/I_{\text{sat}}) - A_{\text{ns}}$$

where *T* is transmittance,  $\Delta T$  is modulation depth, *I* is intensity of laser,  $I_{\text{sat}}$  is saturation power intensity and  $A_{\text{ns}}$  is non-saturable absorbance. Figure 3(a) shows that the modulation depth (MD) of the WS<sub>2</sub>-PVA thin film is 10.17%, the saturating intensity about 8.89 MW cm<sup>-2</sup> and the non-saturable loss about 3.41%. As can be seen from figure 3(b), the Bi<sub>2</sub>Te<sub>3</sub>-PVA thin film exhibits a MD of 23.04%, which is higher than that of the WS<sub>2</sub>-PVA thin film, a saturating intensity of 22.55 MW cm<sup>-2</sup> and a non-saturable loss of 48.39%.

#### 3. Experiment setup

Figure 4 shows the experimental setup for the passively Q-switched Yb-fiber laser based on WS<sub>2</sub>–PVA thin film. The ring laser oscillator cavity comprises a wavelength division multiplexer (WDM), a gain fiber, a polarization-independent isolator (PI-ISO), a polarization controller (PC), a WS<sub>2</sub>–PVA thin film and an optical coupler (OC). A 30 cm long Yb-doped fiber (Liekki Yb 1200-4/125) with an absorption coefficient of 1200 dB m<sup>-1</sup> at 976 nm was used as the gain medium. The 60% port of an optical coupler is used for the laser output. The fiber laser is core pumped with a laser diode emitting at 975.4 nm, with a maximum output power of 700 mW. The PI-ISO was used to force unidirectional operation in the fiber ring cavity. The PC was employed in the cavity to optimize Q-switched operation. The SA film was cut



Figure 2. Raman spectra of (a) WS<sub>2</sub> and (b) Bi<sub>2</sub>Te<sub>3</sub>.



Figure 3. (a) Nonlinear absorption of the WS2-PVA SA. (b) Nonlinear absorption of the Bi2Te3-PVA SA.



Figure 4. Diagram of the experimental setup.

into a small film of  $1 \text{ mm} \times 1 \text{ mm}$ . The small SA film was transferred onto the standard FC/APC fiber end face. The fiber core was entirely covered by the film. The standard FC/APC fiber with the SA film was connected with the other clean one by a standard flange adapter, forming an all-fiber structure. The architecture of the passively *Q*-switched Yb-fiber laser based on the Bi<sub>2</sub>Te<sub>3</sub>–PVA thin film was completely the same; only the SA needs to be changed.

#### 4. Results and discussion

*Q*-switching operation is obtained based on the two types of SA as follows. When the pump power is increased gradually, free

running of the continuous wave (CW) laser is first observed. When the pump power is further increased, *Q*-switched pulses are generated at a certain pump power for different SA materials. The WS<sub>2</sub> *Q*-switching operation starts at a pump power of 50 mW, and Bi<sub>2</sub>Te<sub>3</sub> starts when the pump power is higher than 60 mW. Figure 5 shows the stable *Q*-switching pulse trains of WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> at a certain pump power. Figure 5(a) shows that the repetition rate of the *Q*-switching pulse based on WS<sub>2</sub> is 53.80 kHz at a pump power of 72 mW, corresponding to a pulse interval of 18.58  $\mu$ s. Figure 5(b) shows that the repetition rate of the *Q*-switching pulse based on Bi<sub>2</sub>Te<sub>3</sub> is 75.39 kHz at a pump power of 88 mW, corresponding to a pulse interval of 13.26  $\mu$ s. A larger scanning range of the pulse-train under 100  $\mu$ s is also presented in the insets of figures 5(a) and (b).



Figure 5. Stable *Q*-switching pulse trains of WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> at a certain pump power.



Figure 6. Pulse duration and repetition rate versus incident pump power.



Figure 7. Output power and single pulse energy versus incident pump power.

Figure 6 shows the evolution of the repetition rate and pulse duration with pump power; these are typical features of Q-switching operation. Figure 6(a) shows that the repetition rate of WS<sub>2</sub> pulse trains increases linearly from 36.78 kHz to 81.75 kHz with the increase in pump power, corresponding to the pulse duration reducing from 3.61  $\mu$ s to 1.0  $\mu$ s. Figure 6(b) shows that the repetition rate of Bi<sub>2</sub>Te<sub>3</sub> pulse trains increases linearly from 53.31 kHz to 106 kHz, while the pulse duration varies in the range 1.79  $\mu$ s to 1.07  $\mu$ s. It can be concluded that

 $Bi_2Te_3$  has higher the repetition rate than  $WS_2$  under the maximum pump power, which is an obvious advantage. The shortest pulse duration of the two is very similar. The pulse duration can be further narrowed by shortening the ring laser cavity and increasing the modulation depth of the SA [24–26].

The output power and single pulse energy can be varied by increasing the pump power. This relationship is described in figure 7. The maximum single pulse energy based on  $WS_2$ and  $Bi_2Te_3$  are 56.50 nJ and 61.80 nJ, respectively, under the

**Table 1.** Typical Q-switched results based on WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> in the 1.0  $\mu$ m band in recent reports.

SA	Repetition rate (kHz)	Pulse duration ( $\mu$ s)	Maximum pulse energy (nJ)	Modulation depth (%)	Reference
WS <sub>2</sub>	24.9-36.7	3.2	13.6	3.1	[27]
$WS_2$	60.2–97.0	1.65	28.8	3.87	[19]
$WS_2$	36.78-81.75	3.61-1.0	56.50	10.17	Our work
Bi <sub>2</sub> Te <sub>3</sub>	35–77	1.0-1.3	38.3	2.5	[21]
Bi <sub>2</sub> Te <sub>3</sub>	53.31-106.0	1.79-1.07	61.80	23.04	Our work



**Figure 8.** The dual-wavelength lasing spectrum of the *Q*-switched Yb-doped fiber laser based on  $Bi_2Te_3$  at a pump power of 92 mW. Inset: CW laser output spectrum with a  $Bi_2Te_3$  SA in the cavity.

same pump power. From figure 7(b), we observe that the output power and single pulse energy are reduced under a high incident pump power, but stable Q-switching still occurs. It may be that Bi<sub>2</sub>Te<sub>3</sub> suffers from two-photon absorption under a high incident pump power. In figure 7, when the pump power exceeded 94 mW and 116 mW, for WS<sub>2</sub> and  $Bi_2Te_3$ , respectively, the *Q*-switching state becomes unstable. If the pump power is increased further, the Q-switching state disappears. A stable Q-switching state can be obtained again by decreasing the pump power. The unstable Q-switching phenomenon may be attributed to oversaturation rather than to thermal damage, as typically observed in some other passively Q-switched fiber lasers [23–25]. In table 1, we compare the results for passively Q-switched fiber lasers at the 1.0  $\mu$ m wave band based on WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> SAs. Compared with other reports, it is noted that the maximum single pulse energy and the shortest pulse duration of Bi2Te3 and WS2 are achieved in our work, which is mainly due to the high modulation depth of the SAs.

Finally, when we carefully adjust the PC, stable dualwavelength *Q*-switching operation with  $Bi_2Te_3$  was achieved under a pump power of 92 mW (as observed by the optical spectrum analyzer). As shown in the figure 8, there was dualwavelength simultaneous oscillation at 1026.52 nm and 1029.46 nm. The dual-wavelength output possess a wavelength spacing of 2.94 nm. In order to confirm the dual-wavelength output resulting from the  $Bi_2Te_3$ , we also measured the CW spectrum without the  $Bi_2Te_3$ , which is shown in the inset of figure 8. The center wavelength of the CW spectrum is 1035.33 nm. Because such a dual-wavelength fiber laser has potential applications in microfiber photonics research, microwave carriers and high bit-rate pulses it has gained wide attention [28].

#### 5. Conclusion

In summary, a high-energy Q-switched Yb-doped fiber laser based on WS<sub>2</sub> and Bi<sub>2</sub>Te<sub>3</sub> SAs has been demonstrated. WS<sub>2</sub> shows a modulation depth of 10.17% while Bi<sub>2</sub>Te<sub>3</sub> exhibits a modulation depth of 23.04%. The maximum single pulse energy based on these SAs are 56.50 nJ and 61.80 nJ, respectively. Stable dual-wavelength Q-switching operation was achieved with Bi<sub>2</sub>Te<sub>3</sub>. Comparing the two kinds of SA, Bi<sub>2</sub>Te<sub>3</sub> has obvious advantages.

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