Structural and optical properties of tungsten-doped vanadium dioxide films*

Wang Xue-Jin(王学进)^{a)†}, Liu Yu-Ying(刘玉颖)^{a)}, Li De-Hua(李德华)^{b)}, Feng Bao-Hua(冯宝华)^{b)}, He Zhi-Wei(何志巍)^{a)}, and Qi Zheng(祁 铮)^{a)}

^{a)}College of Science, China Agricultural University, Beijing 100083, China

b) Laboratory of Optical Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China

(Received 3 May 2012; revised manuscript received 26 October 2012)

Thin films of tungsten (W)-doped thermochromic vanadium dioxide (VO_2) were deposited onto soda-lime glass and fused silica by radio frequency magnetron sputtering. The doped VO₂ films were characterized by X-ray diffraction, optical transmittance measurement, and near field optical microscopy with Raman spectroscopy. X-ray diffraction patterns show that the (011) peak of W-doped thermochromic VO₂ film shifts to a lower diffraction angle with the increase of W concentration. The optical measurements indicated that the transmittance change (ΔT) at wavelength of 2500 nm drops from 65% (ΔT at 35 °C and 80 °C for undoped VO₂ film) to 38% (ΔT at 30 °C and 42 °C for the doped VO₂ film). At the same time, phase transition temperature drops from 65 °C to room temperature or lower with the increase of W concentration. Near field optical microscopy image shows that the surface of W-doped VO₂ film is smooth. Raman results show that the main Raman modes of W-doped VO₂ are centered at 614 cm^{-1} , the same as that of undoped VO₂, suggesting no Raman mode changes for lightly W-doped VO₂ at room temperature, due to no phase transition appearing under this condition.

Keywords: tungsten-doped vanadium dioxide, thermochromic, magnetron sputtering DOI: 10.1088/1674-1056/22/6/066803

PACS: 68.55.–a, 81.15.cd, 87.64.km, 87.64.kp

1. Introduction

A thin film of vanadium dioxide (VO_2) undergoes a first order transition from a semiconducting state to a metal state at 68 °C accompanied by the steep change of opto-electronic properties.^[1] Compared with all of the other vanadium oxides, VO₂ thin film has the nearest phase transition temperature (T_t) to room temperature. Moreover, T_t could decrease to or below the room temperature by doping high valance metal ions, such as tungsten, molybdenum, tantalum, etc., into the VO₂ films, which makes it very suitable for fabricating smart windows, infrared (IR) detectors, and high speed switches.^[2-4] VO₂ thin films were deposited by many methods, such as solgel,^[5] sputtering,^[6–9] pulse laser deposition,^[10] atmospheric pressure chemical vapor deposition,^[11] etc. Sol-gel and sputtering are much more efficient to fabricate large-area uniform VO₂ thin films. There are also many methods to dope W into VO₂ thin films partially determined by the deposition methods of VO₂ thin films. [5,11-14] In the case of sputtering, one method is simultaneously sputtering V and W targets into Ar/O2 atmosphere for W-doped VO₂ films,^[14] the other one is preparing W-doped vanadium oxide films by reactive sputtering of a V-W (1.6 at% W) alloy target.^[13] In a dual target configuration, it is easy to control the tungsten concentration in the VO₂ films, however, two dependent systems are needed to control the sputtering of targets, which leads to a high cost. In a V-W (1.6 at% W) alloy target configuration, it is difficult to change tungsten concentraion in VO2 films by tuning sputtering parameters only. If various tungsten concentration V-W targets have to be used, the cost of materials will increase dramtically.

To reduce the cost, we introduce a method to prepare Wdoped VO₂ thin films in this study. Only the V metal (whose diameter is 54 mm) with a small amount of tungsten wires was used as a target to fabricate the uniform W-doped VO₂ films. We found that it was convenient to control the tungsten content in VO₂ films by changing the number of tungsten wires at the glowing zone. This is very important for reducing the cost of materials when large scale production of VO₂ films is carried out in industry.

2. Experimental details

Thin films of VO₂ and W–VO₂ were deposited in an argon/oxygen atmosphere by radio frequency (RF) magnetron sputtering. Before the film deposition, the sputtering chamber was evacuated down to a pressure of about 1×10^{-3} Pa using a turbomolecular pump. Then the argon/oxygen gas with a typical purity of 99.999% was introduced into the sputtering chamber. The vanadium had a purity of 99.9%, tungsten wire on the target, and soda-lime glasses and fused silica with a temperature of about 580 °C were used as the target, the dopant, and substrates, respectively. Typical deposition pa-

*Project supported by the Chinese Universities Scientific Fund (Grant No. 2013QJ007), the Science Fund of China Agricultural University (Grant No. 2007037), the Major Project Foundation of Science and Technology Innovation in Advanced Education (Grant No. 21010112).

[†]Corresponding author. E-mail: xjwang@cau.edu.cn

^{© 2013} Chinese Physical Society and IOP Publishing Ltd

rameters were as follows: the ratio of oxygen to argon was $0.8\% \sim 1.2\%$; the total sputter gas pressure was 1.0 Pa; RF power was 120 W.

Conventional X-ray diffraction (XRD) measurement was performed by means of RINT 2400 with monochromatized Cu K α radiation. The transition spectra of the samples were measured by a spectrometer (BIORAD FTS 6000) in the infrared (IR) region. The Raman spectra were acquired at room temperature using near field optical microscopy with Raman measurement.

Thin films of pure VO₂ and W–VO₂ on various substrates were prepared at the same oxygen partial pressure of 10 mPa while the total pressure of Ar/O₂ is 1 Pa. The W– VO₂ film measured by X-ray photoelectron spectroscopy is $W_{0.014}V_{0.986}O_2$.

3. Results and discussion

The structure and properties of VO₂ films doped with titanium, $Ti_xV_{1-x}O_2$, are stable when *x* changes from 0 to 1,^[15] which is different with that of W-doped VO₂ films. The structure and properties of W-doped VO₂ film are stable if the quantity of tungsten is small. Otherwise, they will be unstable and colored stripes will show up on the sample several days later.

By controlling the content of tungsten carefully, W-doped VO₂ thin films have been successfully deposited onto various substrates. Figure 1 shows the XRD patterns for VO₂ and W-VO₂ films deposited onto soda-lime glass. As shown in Fig. 1, the strong peaks in both patterns correspond to the (011) diffraction of the monoclinic VO₂. The XRD peaks of curve a corresponding to 27.88° , 39.80° , 42.25° , and 57.28° are attributed to the (011), (020), (210), and (022) planes of vanadium dioxide film, respectively. The XRD peaks of curve *b* centered at 27.75°, 37.10°, 39.80°, 55.34°, and 57.28° are attributed to the (011), (200), (020), (220), and (022) planes of vanadium dioxide film, respectively. The inset of Fig. 1 plots the XRD patterns of the above two samples with 2θ ranging from 25° to 31° which shows that the main (011) peak shifts to a lower diffraction angle when the vanadium dioxide film is doped with tungsten and the distance between neighbor (011) planes, d_{011} , increases with the increase of the doping level according to Bragg's law, $d\sin\theta = k\lambda$, where $k = 1, \lambda$ is the wavelength of X-ray, θ is the diffraction angle and d is the distance between the neighbor planes. We suggest that the compressive stress in the (011) plane influences the increase of d_{011} , which will induce a tensile stress along the [011] direction. It is well known that when vanadium dioxide changes from a semiconducting state to a metal state, its volume will expand by 1% of its original volume. Therefore, this tensile stress will be helpful to the phase transition occurring at the lower temperature. We attribute the decrease of T_t of W-doped VO_2 film to the shifting of (011) peaks to the lower diffraction angles.



Fig. 1. (color online) XRD patterns of VO₂ film (curve *a*, red) and W–VO₂ film (curve *b*, blank). The inset is the XRD patterns with 2θ ranging from 25° to 31°.

The transmittance for W-doped VO2 films versus the wavelength at different test temperatures are shown in Fig. 2. It can be seen that the transmittance in the whole wavelength changes little at temperatures ranging from 46 °C to 42 °C, the transmittance changes with temperature at a rate of about 1%/°C at the wavelength of 2500 nm. However, the transmittance changes obviously at the temperature ranging from 42 °C to 40 °C, its change rate increases clearly up to 2%/°C at 2500 nm. The transmittance changes rapidly at 2500 nm at the temperature ranging from 40 °C to 30 °C, the transmittance changes with temperature according to 3%/ °C. The transmittance as a function of the temperature at the wavelength of 2500 nm for W-doped has been drawn in Fig. 3. For comparison, the temperature-depend transmittance at 2500 nm for undoped VO₂ films deposited on fused silica and soda-lime glass has also been plotted. The phase transition temperature $T_{\rm t}$ of W–VO₂ film deposited on soda-lime glass is about 35 °C in Fig. 2, and the T_t of VO₂ films deposited on soda-lime glass and fused silica are 63 °C and 69 °C, respectively. The reason for the changing of T_t of VO₂ films on various substrates can be found in our previous paper.^[16]



Fig. 2. (color online) Transmittance as a function of wavelength at different temperatures for W–VO₂ film.

Raman spectra were used to confirm the presence of monoclinic VO₂, as well as W-doped VO₂.^[17–19] Piccirillo *et al.* reported that the Raman spectrum of doped VO₂ film, such as $V_{0.9825}W_{0.0175}O_2$, is the same as the spectrum assigned to the monoclinic VO₂, and the spectrum also shows no sign of carbon contamination determined by the peak at 1355 cm^{-1} in the films.^[17] From the curves a and b of Fig. 4, it can be seen that all the signals detected and marked in the spectra are the same, and the carbon contamination in both films is not obvious. The major peaks of the spectra of both films are centered at 223, 260, 303, 392, and 614 cm^{-1} , respectively, which is in agreement with the literature data reported elsewhere.^[19] No other vanadium oxide is present and the Raman spectrum of monoclinic VO_2 is almost the same as that of W-doped VO_2 . We point out that impurities such as W ions instead of V ions do not change the structure of VO2 significantly. Although the ratio of vanadium to oxygen slightly deviates, the stoichiometric VO₂ would dramatically influence the peak positions of Raman spectra. This is probably due to the W ions instead of V ions located in the lattice spots of monoclinic VO2 which, despite a slight distortion, does not mechanically breakdown the structure, so that the vibration modes of both W-doped VO₂ and undoped VO₂ remain unchanged.



Fig. 3. Transmittance at wavelength 2500 nm versus temperature for W–VO₂ (phase transition temperature $T_{\rm t} = 35$ °C) and undoped VO₂ ($T_{\rm t} = 63$ °C) films on the soda-lime glass, and for undoped VO₂ film ($T_{\rm t} = 69.5$ °C) on the fused silica.



Fig. 4. (color online) Raman spectra of VO₂ film (curve a, red) and W–VO₂ film (curve b, blank).

4. Conclusion

We present a study on the structural and optical properties of pure and W-doped vanadium dioxide films. We have shown that an increase in tungsten content in the film results in a slight shift of the diffraction angle towards a lower angle.

Upon going through semiconductor-metal transition, the transmittance of the doped VO₂ film does not drop as dramatically as that of the pure VO₂ film. The transmittance change of the doped VO₂ film is about 38% (ΔT at 30 °C and 42 °C) at 2500 nm, while that of the pure VO₂ film is about 65% (ΔT at 35°C and 80°C).

We have shown that Raman spectra of both doped and undoped VO₂ films are nearly the same. We suggest that although W ions alternating V ions located in the lattice spots of monoclinic VO₂ distort the lattice slightly, they do not induce phase transition. Therefore, the vibration modes of W-doped VO₂ and undoped VO₂ are the same.

Acknowledgements

The authors thank professor Feng Ke-An (Institute of Physics, Chinese Academy of Sciences) for the fruitful discussions, and Dr. Xu Jin-Jie for the Raman spectra measurement.

References

- [1] Kivaisi R T and Samiji M 1999 Sol. Energy Mater. Sol. Cells 57 141
- [2] Granqvist C G 2007 Sol. Energy Mater. Sol. Cells 91 1529
- [3] Wang H C, Yi X J, Chen S H, Huang G and Xiao J 2004 J. Infrared Millim. Wave 23 99 (in Chinese)
- [4] Wang H C, Yi X J, Li Y 2005 Opt. Commun. 256 305
- [5] Chae B G, Kim H T, Yun S J, Kim B J, Lee Y W, Youn D H and Kang K Y 2006 *Electrochem. Solid-State Lett.* 9 C12
- [6] Jin P and Tanemura S 1994 Jpn. J. Appl. Phys. Part 1 33 1478
- [7] Lee M H and Cho J S 2000 Thin Solid Films 365 5
- [8] Wang X J, Liang C J, Guan K P, Li D H, Nie Y X, Zhu S Q, Huang F, Zhang W W and Cheng Z W 2008 *Chin. Phys. B* 17 3512
- [9] Luo Z F, Wu Z M, Xu X D, Wang T and Jiang Y D 2010 Chin. Phys. B 19 106103
- [10] Soltani M, Chaker M, Haddad E and Kruzelesky R V 2006 J. Vac. Sci. Technol. A 24 612
- [11] Manning T D and Parkin I P 2004 J. Mater. Chem. 14 2554
- [12] Jin P, Nakao S and Tanemura S 1998 Nucl. Instrum. Method Phys. Res. B 141 419
- [13] Jin P, Tazawa M, Ikeyyama M, Tanemura S, Macak K, Wang X, Olafsson S and Helmersson U 1999 J. Vac. Sci. Technol. A 17 1817
- [14] Romanyuk A, Steiner R, Marot L and Oelhafen P 2007 Solar Energy Material and Solar Cell 91 1831
- [15] Kakiuchida H, Jin P and Tazawa M 2008 Thin Solid Films 516 4563
- [16] Wang X J, Xu J J, Fei Y J, Li D H, Li T J, Nie Y X, Feng K A and Wu N J 2002 Jpn. J. Appl. Phys. Part 1 41 312
- [17] Piccirillo C, Binions R and Parkin I P 2007 Chem. Vap. Deposition 13 145
- [18] Piccirillo C, Binions R and Parkin I P 2008 Thin Solid Films 516 1992
- [19] Parker J C 1990 Phys. Rev. B 42 3164